

## Lecture 6: Criteria for QECC

6.1 a simple classical ECC:

repetition code

$$\text{encoding: } C: \{0,1\} \rightarrow \{0,1\}^3$$

$$x \mapsto \underset{x}{\underbrace{xxx}}$$

send every bit (symbol) three times (or  $n$  times)

Error correction step:

- compare all received symbols  $\vec{y} = (y_1, \dots, y_n)$
- use the most frequent as the decoding result

for  $n=3$ :

1 error can be corrected

$n=2t+1$

$t$  errors can be corrected

Assumption:

a small number errors is more likely than many errors

rate of the encoding:  $\frac{1}{n}$

number of errors per symbol  $\frac{t}{2t+1} \rightarrow \frac{1}{2}$

A direct analogue of the repetition code does not work for quantum information since

1. quantum quantum states cannot be replicated
2. multiple copies of a quantum state, some of which might be corrupted, cannot be compared (without destroying them)

### 6.2 A first quantum code

bad idea: encoding  $| \phi \rangle \rightarrow | \phi \rangle^{\otimes n}$

good idea: take copies of the basis states:

$$\text{C: } \mathbb{C}^2 \rightarrow (\mathbb{C}^2)^{\otimes 3}$$

$$| 0 \rangle \mapsto | 000 \rangle$$

$$| 1 \rangle \mapsto | 111 \rangle$$

$$| \phi \rangle = \alpha | 0 \rangle + \beta | 1 \rangle \mapsto \alpha | 000 \rangle + \beta | 111 \rangle$$

(3)

|  | state   | error operator                 | space           |
|--|---|--------------------------------|-----------------|
| no error                               | $\alpha 000\rangle + \beta 111\rangle$          | $I \otimes I \otimes I$        | $\mathcal{C}_0$ |
| $\sigma_x$ on 1 <sup>st</sup> position | $\alpha 100\rangle + \beta 011\rangle$          | $\sigma_x \otimes I \otimes I$ | $\mathcal{C}_1$ |
| $\sigma_x$ on 2 <sup>nd</sup> position | $\alpha 0\cancel{0}\rangle + \beta 101\rangle$  | $I \otimes \sigma_x \otimes I$ | $\mathcal{C}_2$ |
| $\sigma_x$ on 3 <sup>rd</sup> position | $\alpha 001\rangle + \beta 11\cancel{0}\rangle$ | $I \otimes I \otimes \sigma_x$ | $\mathcal{C}_3$ |

$$\mathcal{C}_0 = \langle |000\rangle, |111\rangle \rangle$$

$$\mathcal{C}_1 = \langle |100\rangle, |011\rangle \rangle = (\sigma_x \otimes I \otimes I) \mathcal{C}_0$$

$$\mathcal{C}_2 = \langle |0\cancel{0}\rangle, |101\rangle \rangle = (I \otimes \sigma_x \otimes I) \mathcal{C}_0$$

$$\mathcal{C}_3 = \langle |001\rangle, |11\cancel{0}\rangle \rangle = (I \otimes I \otimes \sigma_x) \mathcal{C}_0$$

$$(\mathbb{C}^2)^{\otimes 3} = \mathcal{C}_0 \oplus \mathcal{C}_1 \oplus \mathcal{C}_2 \oplus \mathcal{C}_3$$

decomposition of  $(\mathbb{C}^2)^{\otimes 3}$  into 4 mutually orthogonal subspaces  $\mathcal{C}_i$

This code can correct one "bit-flip error"  $\sigma_x$ ,  
but cannot even detect a single "sign-flip error"  $\sigma_z$ .

A code correcting one  $\delta_z$ -error (but no  $\delta_x$ -error). (7)

The Hadamard transform  $H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$   
 interchanges  $\delta_x$  and  $\delta_z$ , i.e.

$$H \delta_x H = \delta_z$$

Using the code  $(H \otimes H \otimes H) \rho_0$ , we can  
 correct one sign-flip error

$$(H \otimes H \otimes H) |000\rangle = \frac{1}{\sqrt{8}} \sum_{x \in \{0,1\}^3} |x\rangle$$

$$(H \otimes H \otimes H) |111\rangle = \frac{1}{\sqrt{8}} \sum_{x \in \{0,1\}^3} (-1)^{x_1+x_2+x_3} |x_1 x_2 x_3\rangle$$

different bases

$$|0\rangle \rightarrow \sum_{\text{wt}(x) \text{ even}} |x\rangle$$

$$|1\rangle \rightarrow \sum_{\text{wt}(x) \text{ odd}} |x\rangle$$

$\text{wt}(x)$  = number of non-zero elements

(5)

first "full" quantum error-correcting code (QECC)  
by Peter Shor:

Combine two levels of these simple codes,  
one correcting  $\delta_x$ -errors, the other level  
correcting  $\delta_z$ -errors

$\Rightarrow$  nine qubit code correcting one "general"  
error (more details later)

### 6.3 QECC conditions by Knill & Laflamme 97

(6)

Let  $\mathcal{C} \subseteq \mathcal{H}$  be a subspace of the Hilbert space  $\mathcal{H}$  and let  $\{A_i : i \in I_Q\}$  be the errors of a quantum channel  $Q$  on  $\mathcal{H}$ .

Furthermore, let  $\{|c_i\rangle : i \in I_E\}$  be an orthonormal basis of  $\mathcal{C}$ .

Then  $\mathcal{C}$  is a quantum error-correcting code for the channel  $Q$ , iff

- (i)  $\forall k, l \in I_Q \forall i \neq j \in I_E : \langle c_i | A_k^+ A_l | c_j \rangle = 0$
- (ii)  $\forall k, l \in I_Q \forall i, j \in I_E : \langle c_i | A_k^+ A_l | c_i \rangle = \langle c_j | A_k^+ A_l | c_j \rangle = \alpha_{kl}$

Equivalently, if  $P_E$  denotes the projection onto  $\mathcal{C}$ , then

- (iii)  $\forall k, l \in I_Q : P_E A_k^+ A_l P_E = \alpha_{kl} P_E$

## 6.4 Lemma:

The error-correction conditions 6.3 are linear in the error operators.

Proof:  $A = \sum \lambda_k A_k \quad B = \sum \mu_e A_e$

a)  $\langle c_i | A^+ B | c_j \rangle = \langle c_i | (\sum \lambda_k^* A_k^*) (\sum \mu_e A_e) | c_j \rangle$   
 $= \sum \lambda_k^* \mu_e \langle c_i | A_k^* A_e | c_j \rangle \stackrel{(ii)}{=} 0$

b) similarly:  $\langle c_i | A^+ B | c_i \rangle \stackrel{(ii)}{=} \langle c_i | A^+ B | c_j \rangle$

Consequence: It is sufficient to test / fulfill the conditions 6.3 for a vector space basis of all error operators of the channel.

# "constructive" proof of the QECC conditions 6.3

in the condition, we have inner products of the (unnormalized) states  $A_k |c_i\rangle$

|                 | $A_1 \mathcal{C}$ | $A_2 \mathcal{C}$ | ... | $A_k \mathcal{C}$ |
|-----------------|-------------------|-------------------|-----|-------------------|
| $\mathcal{V}_0$ | $A_1  c_0\rangle$ | $A_2  c_0\rangle$ | ... | $A_k  c_0\rangle$ |
| $\mathcal{V}_1$ | $A_1  c_1\rangle$ | $A_2  c_1\rangle$ | ... | $A_k  c_1\rangle$ |
|                 | :                 |                   |     |                   |
| $\mathcal{V}_i$ | $A_1  c_i\rangle$ | $A_2  c_i\rangle$ | ... | $A_k  c_i\rangle$ |
|                 |                   |                   | ... |                   |

Condition (ii): the spaces  $\mathcal{V}_i$  are mutually orthogonal,  
i.e  $\mathcal{V}_i \perp \mathcal{V}_j$

Consider the vector space  $\mathcal{V}_0 = \{A_k |c_0\rangle : k \in \mathbb{I}_Q\}$   
perform a Gram-Schmidt orthogonalization in  $\mathcal{V}_0$

for  $j \in I_Q$  do

$$|b_j^{(i)}\rangle \leftarrow A_j |c_i\rangle$$

$$|b_j^{(i)}\rangle \leftarrow |b_j^{(i)}\rangle - \sum_{l < j} \langle b_l^{(i)} | b_j^{(i)} \rangle \cdot |b_l^{(i)}\rangle$$

$$\text{if } |b_j^{(i)}\rangle \neq 0$$

$$|b_j^{(i)}\rangle \leftarrow \frac{1}{\sqrt{\langle b_j^{(i)} | b_j^{(i)} \rangle}} \cdot |b_j^{(i)}\rangle$$

end if

end for

$\Rightarrow$  ONB for the space  $V_i$

$$|b_j^{(i)}\rangle = \sum_{k \in I_Q} d_{jk}^{(i)} A_k |c_i\rangle$$

The coefficients  $d_{jk}^{(i)}$  depend only on the inner products of the vectors  $A_k |c_i\rangle$  in  $V_i$ , i.e. on

$\Rightarrow$  They are independent of the state  $|c_i\rangle$ .

Instead of taking linear combinations of the basis states in  $\mathcal{V}_i$ , we replace the operators  $A_k$  by linear combinations, i.e.

$$\tilde{A}_j := \sum_{k \in I_Q} \lambda_{jk} A_k$$

$\Rightarrow$  new scheme:

|                 | $\tilde{A}_1 \mathcal{C}$ | $\tilde{A}_2 \mathcal{C}$ | ... | $\tilde{A}_n \mathcal{C}$ |
|-----------------|---------------------------|---------------------------|-----|---------------------------|
| $\mathcal{V}_0$ | $\tilde{A}_1  c_0\rangle$ | $\tilde{A}_2  c_0\rangle$ | ... | $\tilde{A}_n  c_0\rangle$ |
| $\mathcal{V}_1$ | $\tilde{A}_1  c_1\rangle$ | $\tilde{A}_2  c_1\rangle$ |     | $\tilde{A}_n  c_1\rangle$ |
| $\vdots$        |                           |                           |     |                           |
| $\mathcal{V}_i$ | $\tilde{A}_1  c_i\rangle$ | $\tilde{A}_2  c_i\rangle$ |     | $\tilde{A}_n  c_i\rangle$ |

$\Rightarrow$  now the spaces  $\tilde{A}_k \mathcal{C}$  and  $\tilde{A}_e \mathcal{C}$  for  $k \neq l$  are mutually orthogonal

$\rightarrow$  There exists a measurement which projects onto one of these spaces and yields information on the error  $\tilde{A}_k$ .